



Both micro-structural slips and conductivity variation properties of magneto nanoflow of micropolar fluid



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Abstract

Electrical and thermal conductivity plays a paramount role in food industries' heat operation processes. Hence, this analysis sheds light on the variation of conductivity properties in MHD boundary layer flow with micropolar nanofluid. Micro-structural slips and Arrhenius impacts are also taken into consideration. Reliable suitable similarity transformations of the resulting fluid equation system are applied, to give an ordinary differential system of equations. Rung-Kutta method with shooting technique is utilized to achieve numerical results. The technical algorithms are made via a full package of *Mathematica 13.1.1*. Sketches/results are prepared for the velocity, microrotation velocity, temperature and nanoparticles concentration distributions against physical parameters. Results assured on the variation of conductivity properties improve the temperature distribution of fluid and then the operation system attained more energy, that beneficial in industries of food operations with the aid idea of Ohmic heaters.

Keywords: Electrical and thermal conductivity in variable case; Slip velocity in a multiple case; Micropolar fluid; Mathematica package ver. 13.

1. Introduction

Nowadays, in many industrial processes the electrical conductivity plays the most important rule, which considered phenomenon that happens in some materials once cooled to temperatures in lower degree. In 1911 [1] Onnes offers the first look at the meaning of electrical conductivity and obtain the definite resistance of metals at small temperatures. In the latest view, scholars, and examiners premeditated the electrical conductivity bases to its numerous claims: In electrical processes such magnetic resonance imaging devices, digital electrical circuits, particle accelerators, in medicine processes, like magnetic resonance. Inflows, Obalalu *et al* [2] announced a recommendation for the electrical conductivity variable case influence of Casson fluid in presence of nanoparticle. They alleged that electrical conductivity has a converse relation with fluid concentration. Otherwise, Qasim *et al* [3] suggest a new mechanism of nanoparticle theory on the peristaltic pumping. Electrical conductivity has a dominant relation to fluid concentration and

temperature. An original theoretical method to the relationship among electrical conductivity and fluid concentration. The literature on the line of conductivity properties and their relation to fluid concentration is in Refs. [4 – 8].

The flows with slip conditions are considered a substantial concern of microsystems in applications like micro-valves, pumps with small sizes, and nozzles. The boundary layer flow with slip conditions in the case of multiple case was constructed by Dawar *et al* [9]. Authors show that microstructural slip parameter is the cause growing in velocity distribution more than in the case of a non-microstructural slip. Velocity slip and thermal radiation effects on the unsteady flow with stretching sheet problem are numerically solved by Mabood and Shateyi [10]. The author state that the slip parameter is the cause in enhancing the fluid velocity at the center of the layer and diminish the velocity at the edge of layer. Mahmoud and Waheed [11] scrutinized the influence of the magnetic field on micropolar fluid with boundary layer flow. Afzal and Aziz [12] addressed

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the effect of MHD slip and non-constant conductivity properties on the flow of nanofluid in the flow with presence of nanoparticles. In addition, many researches are studied the important/ substantial applications of fluid in different types(see Refs. [13 – 20]).

The results of such fluid problems need a high technical algorithm to cure the high nonlinear terms that appear in MHD fluid of the boundary flow, So, we select a semi-analytical method which named by algorithms of the shooting method joint with the iteration of Rung-Kutta-Fehlberg method. **Generally**, the shooting algorithm is measured as an extremely helpful method that is secondhand to solve the problems of boundary value by reducing it to 1st problem. The algorithm of the shooting algorithms is anticipated initially by Merle 1988 [21], who finds the solution of nonlinear Dirac equations. The shooting technique is discretized for solving flow studies with transpiration influence by Hasanuzzaman et al [22]. Lanjwani et al [23] studied the Casson fluid flow using the shooting method. Many authors and detectives use the numerical and semi-analytical algorithms to get the solutions of boundary layer problems with variable conditions, see Refs. [24 – 40].

In the above investigations, all researchs are in the line of boundary layer flow problems with constant conductivity properties. This is not sufficient to visualize all their applications in food industries. The existing paper presents a new theoretical suggestion for the mechanisms of conductivity properties (Thermal and electrical) varying with the fluid concentration. This suggestion helps to provide the behavior of non-constant electrical and thermal conductivity in many industries' processes. Sketches are extracted for all possible distributions of velocities and fluid temperature and concentration. This study can serve as a model which may support in comprehension the mechanics of physiological flows. Physically, our model corresponds to transport of the gastric juice in the small intestine.

2 Formulation of the problem

The flow of a micropolar nanofluid on a surface with multiple slips and stretches. Non-constant electrical and thermal conductivity is premeditated. $\check{u}_s = a\check{x}$ is supposed the stretching velocity and a is the rate of stretch. The fluid equation model is theoretically scrutinized as the following [9]:

$$\frac{\partial \check{u}}{\partial \check{x}} + \frac{\partial \check{v}}{\partial \check{y}} = 0, \tag{1}$$

$$\check{u} \frac{\partial \check{u}}{\partial \check{x}} + \check{v} \frac{\partial \check{u}}{\partial \check{y}} = \frac{\kappa}{\rho} \frac{\partial \check{N}}{\partial \check{y}} + \frac{(\mu + \kappa)}{\rho} \frac{\partial^2 \check{u}}{\partial \check{y}^2} - \frac{\sigma(\check{C})B_0^2 \check{u}}{\rho} + g\{\beta_c(\check{C} - \check{C}_\infty) + \beta_T(\check{T} - \check{T}_\infty)\}, \tag{2}$$

$$\check{u} \frac{\partial \check{N}}{\partial \check{x}} + \check{v} \frac{\partial \check{N}}{\partial \check{y}} = \frac{\Omega}{\rho_j} \frac{\partial^2 \check{N}}{\partial \check{y}^2} - \frac{\kappa}{\rho_j} \left(2\check{N} + \frac{\partial \check{u}}{\partial \check{y}} \right), \tag{3}$$

$$\check{u} \frac{\partial \check{T}}{\partial \check{x}} + \check{v} \frac{\partial \check{T}}{\partial \check{y}} = \frac{1}{\rho c_p} \left(\frac{\partial}{\partial \check{x}} k(\check{C}) \frac{\partial \check{T}}{\partial \check{x}} + \frac{\partial}{\partial \check{y}} k(\check{C}) \frac{\partial \check{T}}{\partial \check{y}} \right) + \frac{16\sigma^* \check{T}_\infty^3}{3\rho c_p k^*} \frac{\partial^2 \check{T}}{\partial \check{y}^2} + \frac{\sigma(\check{C})B_0^2 \check{u}^2}{\rho c_p} + \check{\tau} \left[D_B \frac{\partial \check{T}}{\partial \check{y}} \frac{\partial \check{C}}{\partial \check{y}} + \frac{D_B}{D_T} \left(\frac{\partial \check{T}}{\partial \check{y}} \right)^2 \right], \tag{4}$$

$$\check{u} \frac{\partial \check{C}}{\partial \check{x}} + \check{v} \frac{\partial \check{C}}{\partial \check{y}} = -K_r^2 (\check{C} - \check{C}_\infty) \left(\frac{\check{T}}{\check{T}_\infty} \right)^n e^{\left(\frac{-E_a}{K_B \check{T}} \right)} + \frac{D_T}{\check{T}_\infty} \frac{\partial^2 \check{T}}{\partial \check{y}^2} + D_B \frac{\partial^2 \check{C}}{\partial \check{y}^2}. \tag{5}$$

The gradient velocity of spin is suited as:

$$\Omega = \mu \left(\frac{\alpha}{2} + 1 \right) j. \tag{6}$$

The parameter α is defined as a micropolar parameter, which equal $(= \frac{\kappa}{\mu})$

The conditions of flow problem are as follows:

$$\check{u} = \check{u}_s + \check{u}_{slip}, \check{v} = 0, \check{T} = \check{T}_s, \check{N} = -\check{n} \frac{\partial \check{u}}{\partial \check{y}}, \check{C} = \check{C}_s \text{ at } \check{y} = 0, \tag{7}$$

$$\check{u} \rightarrow 0, \check{N} \rightarrow 0, \check{T} \rightarrow \infty, \check{C} \rightarrow \infty \text{ as } \check{y} \rightarrow \infty, \tag{8}$$

Now,

$$\check{u}_{slip} = -\frac{2}{3} \left\{ \frac{3}{8} \left[l^4 + \frac{2(1-l^2)}{k_n^2} \right] \lambda^2 \frac{\partial}{\partial \check{y}} \left(-R_0 \check{N} + \frac{\partial \check{u}}{\partial \check{y}} \right) - \lambda \left(\frac{1-\beta l^2}{\beta} - \frac{3}{2} \frac{1-l^2}{k_n} \right) \frac{\partial \check{u}}{\partial \check{y}} \right\}, \tag{9}$$

$$\check{u}_{slip} = \frac{\partial}{\partial \check{y}} \left(-R_0 \check{N} + \frac{\partial \check{u}}{\partial \check{y}} \right) B + \frac{\partial \check{u}}{\partial \check{y}} A, \tag{10}$$

The symbols $A, B, \check{n}, \check{u}, \check{x}, \check{y}, R_0, B_0$ represent the slip constants, the parameter of micro-rotation, the coordinates, the multiple slips, the magnetic field forte respectively, and the concentration (\check{C}) of fluid, the Brownian (D_B) diffusion coefficient, the diffusion (D_T) coefficient, gravitational (g) acceleration, mean absorption coefficient, the Stefan-Boltzmann (σ^*) constant, the Knudsen (k^*) number, the temperature, the expansion (β_c) of concentration, thermal (β_T) expansion, the vortex viscosity, the electrical (κ) conductivity, the density (ρ), the heat ratio, and the dynamic (μ) viscosity, respectively. The similarity transformations are defined as:

$$\check{u} = \frac{\partial \psi}{\partial \check{y}}, \psi = \sqrt{a v x} f(\xi), \check{v} = -\frac{\partial \psi}{\partial \check{y}}, \check{N} = a x \sqrt{\frac{a}{v}} g(\eta), \Theta(\eta) = \frac{\check{T} - \check{T}_\infty}{\check{T}_s - \check{T}_\infty}, \Phi(\eta) = \frac{\check{C} - \check{C}_\infty}{\check{C}_s - \check{C}_\infty}, D_1 = \frac{\check{T}_s - \check{T}_\infty}{x}, \eta = \sqrt{\frac{a}{v}} y, D_2 = \frac{\check{C}_s - \check{C}_\infty}{x}, \tag{11}$$

In the last decade, a researcher [29] present a new relation for conductivity parameters depending on the fluid concentration.

$$\kappa(\hat{C}) = \kappa_0 \left[1 + \beta_1 \left(\frac{\hat{C} - \hat{C}_\infty}{\hat{C}_s - \hat{C}_\infty} \right) \right], \quad (12)$$

Here, the mean thermal (κ_0) conductivity of the suspension, the relation for the dependence $\beta_1 = \frac{1}{\kappa_0} \frac{\partial \kappa}{\partial \hat{C}}$ of the thermal conductivity on the concentration of the nanoparticles, in addition, $0 < \beta_1 < 1$, $\beta_1 = 0$ at the conductivity constant case. The conductivity properties in the conductivity case are known as:

$$\sigma(\hat{C}) = \sigma_0 \left[1 + \beta_2 \left(\frac{\hat{C} - \hat{C}_\infty}{\hat{C}_s - \hat{C}_\infty} \right) + \dots + O(\alpha) \right], \quad (13)$$

As the high order of α leads to minimum values of $\sigma(\hat{C})$ so, we neglect the higher order of conductivity as:

$$\sigma(\hat{C}) = \sigma_0 \left[1 + \beta_2 \left(\frac{\hat{C} - \hat{C}_\infty}{\hat{C}_s - \hat{C}_\infty} \right) \right], \quad (14)$$

Using Eq. (9) - (14), the Eqs. (2)– (5) are condensed as:

$$\gamma_2 \varphi + (\alpha + 1) f'''' - f'^2 - \alpha g' + f f'' + \gamma_1 \theta - M(1 + \beta_1 \theta) f' = 0, \quad (15)$$

$$\left(\frac{\alpha}{2} + 1 \right) g'' + f g' - g f' - \alpha (f'' + 2g) = 0, \quad (16)$$

$$\frac{(1+R)}{Pr} \theta'' + (\beta_2 \theta' \varphi' + (1 + \beta_2 \varphi) \theta'') + E_c M f'^2 + Nb \theta' \varphi' + Nt \theta'^2 - \theta f' + f \theta' + M(1 + \beta_1 \theta) E_c f'^2 = 0, \quad (17)$$

$$\varphi'' - S_c f' \varphi + S_c f \varphi' + \frac{Nt}{Nb} \theta'' - Pr \sigma (1 + \delta \theta)^n e^{\left(\frac{-E}{1 + \delta \theta} \right)} \varphi = 0. \quad (18)$$

The similarity of boundary conditions

$$f'(0) = 1 + \gamma f''(0) + \delta_1 f'''(0) + R_0 g'(0), \quad f(0) = 0, \quad g(0) = -m f'', \quad \theta(0) = 1, \quad \phi(0) = 1, \quad (19)$$

$$\theta(\infty) = 0, \quad \phi(\infty) = 0, \quad f'(\infty) = 0, \quad g(\infty) = 0. \quad (20)$$

Where, the magnetic ($M = \frac{\sigma_0 B_0^2}{\rho \alpha}$) field, the thermal ($R = \frac{16 \sigma^* \bar{T}_\infty^3}{3 \rho c_p k^*}$) radiation parameter, the Soret ($S_c = \frac{\nu}{D_B}$)

Parameter, the thermophoresis ($N_t = \frac{\tau D_T (\bar{T}_s - \bar{T}_\infty)}{\bar{T}_\infty \nu}$) parameter, the Prandtl ($Pr = \frac{\mu c_p}{k}$) parameter, the Brownian motion ($N_b = \frac{\tau D_B (\hat{C}_s - \hat{C}_\infty)}{\bar{T}_\infty \nu}$) parameter, the Eckert ($E_c = \frac{\bar{u}_s^2}{c_p (\bar{T}_s - \bar{T}_\infty)}$) number, the thermal Grashof ($\gamma_1 = \frac{\beta_T D_1 g}{a^2}$) number, the mass Grashof ($\gamma_2 = \frac{\beta_C D_2 g}{a^2}$) number, the chemical ($\sigma = \frac{k_T^2}{a}$) reaction parameter, and the activation ($E = \frac{E_a}{K_B \bar{T}}$) parameter.

3 Adaptive solutions using Shooting

The solution of a system (15-18) with related conditions (19-20) using Rung-Kutta method with the adaptive shooting method is through this section. 1st order of solution for equations is attained via 1st iteration, and then the shooting algorithm is used to 2st, 3st, etc.... If we suppose,

$$f = y_1, \quad f' = y_2, \quad f'' = y_3, \quad g = y_4, \quad g' = y_5, \quad \theta = y_6, \quad \theta' = y_7, \quad \varphi = y_8 \quad \text{and} \quad \varphi' = y_9,$$

Then the recurrence relations for fluid distributions are as:

$$y_3' = \frac{M(1 + \beta_1 y_8) f' [\eta] - \gamma_1 y_6 - \gamma_2 y_8 + y_2^2 + \alpha y_5 - \gamma_1 y_3}{(1 + \alpha)}, \quad (21)$$

$$y_5' = \frac{y_4 y_2 - y_1 y_5 + \alpha (2y_1 + y_3)}{(1 + \frac{\alpha}{2})}, \quad (22)$$

$$y_7' = -\frac{(1+R)(\beta_2 y_7 y_9)}{Pr} + Pr \frac{y_6 y_2 - e c M (1 + \beta_1 y_8) y_2^2 - y_1 y_7 - N t y_7^2 - N b y_7 y_9}{(1 + \beta_2 y_8)}, \quad (23)$$

$$y_9' = -e \frac{e a}{1 + \xi y_6} pr \sigma (1 + \xi y_6) y_8 + y_1 y_9 + \frac{N t y_7'}{N b}, \quad (24)$$

Subject to conditions

$$y_2(0) = 1 + \gamma y_3(0) + \delta_1 y_3'(0) + R_0 y_5(0), \quad y_1(0) = 0, \quad y_4(0) = -m y_3,$$

$$y_6(0) = 1, \quad y_8(0) = 1 \quad (25)$$

$$y_6(\infty) = 0, \quad y_8(\infty) = 0, \quad y_2(\infty) = 0, \quad y_4(\infty) = 0. \quad (26)$$

Therefore, The process is repeated iteratively many times until convergence is occurred, i.e., until the absolute values of the difference between every two successive approximations of the missing conditions is less than (in our problem ϵ is taken $= 10^{-7}$, and these algorithms of the shooting method are used with aid of **Mathematica 13.1.1**.to get the results offered in the next section.

4 Results and discussions

In this section, we sketch of velocity $f'(\eta)$, microrotation velocity $g(\eta)$, temperature $\theta(\eta)$, and nanoparticles concentration $\varphi(\eta)$ profiles. All graphs are offered against pertinent physical parameters of interest for four different values. The following values are used as standard values of the problem parameters: $Pr = 1, Nb = 0.2, M = 0.5, \gamma = 0.1, \alpha = 1, \gamma_2 = 0.6, \xi = 0.2, \beta_1 = 0.2, \beta_2 = 0.2, \sigma = 0.2, E_a = 1, \gamma_1 = 0.5, R = 1, Nt = 0.3, ec = 0.15, \delta = 0.01, b = 0.1, m = 0.1$.

Mass Grashof number is a non-dimensional parameter and is used in the correlation of mass transfer due to the attractive force between particles induced mixed convection at a solid surface immersed in a fluid. The effects of mass Grashof number γ_2 and the magnetic field parameter M on the velocity $f'(\eta)$ are shown in **Figs. (1)** and **(2)**, respectively, and it is shown from these figures that the velocity increases by increasing γ_2 , while it decreases as M increases. Also, for small values of γ_2 and large values of M , the velocity decreases with η , with a relationship that seems like a parabola. The physical expectation and previous definition of mass Grashof number are consistent with the result in Fig. (1). The following the result in figure (2); the force exerted on a charged particle entire a liquid which is moving with a specific velocity through an electric field and magnetic field is called Lorentz force. This entire electromagnetic force on this liquid always prevents the flow. Fig. (3) shows the variation of the velocity $f'(\eta)$ with η for various values of first order slip condition parameter γ . It is clear from Fig.(3), that the velocity decreases as γ increases in the interval $\eta \in [0, 1.78]$; otherwise, it decreases by increasing γ . Therefore, the behavior of $f'(\eta)$ in the interval $\eta \in [1.78, \infty)$ is opposite to its behavior in the interval $\eta \in [0, 1.78]$. The effects of other parameters are similar to those obtained in Figs. (1) and (2).

Eq. (16) evaluates how the microrotation velocity distribution $g(\eta)$ variations with the dimensionless coordinate η . The effects of the microrotation parameter α and m on the microrotation velocity distribution $g(\eta)$ are illustrated in figures (4) and (5), respectively. It is obviously that the microrotation velocity distribution increases as both α and m increase, but in case of m , it seems like a parabola with upper and down vortices for $m < 0.3$, and $m > 0.3$. moreover, the microrotation velocity has a positive sign, and for major values of α and little values of m , it becomes larger with η till a finite value of η , after this maximum value it goes lower. The influence of

thermal Grashof number γ_1 on the microrotation velocity distribution $g(\eta)$ is shown in Fig. (6). It is observed that the microrotation velocity diminutions as γ_1 in the range of η , namely, $\eta \in [0, 1.41]$; otherwise there is an enhancement in the value of $g(\eta)$ as γ_1 increases. So, the manner of $g(\eta)$ in the interval $\eta \in [1.41, \infty)$, is opposite to its manner in the interval $\eta \in [0, 1.41]$, and in the first interval, a maximum value of $g(\eta)$ occurs at $\eta = 0.58$.

The temperature distribution θ is elucidated for various values of the parameter of thermal radiation R and the parameter of microrotation α , is elucidated in Figs. (7) and (8), respectively. It is notted from figures (7) and (8), respectively, that the temperature becomes larger with the increase of R , while it goes lower as α increases. Furthermore, It is observed that for major values of R and little values of α , the relation between θ and η is a straight line, and all curves intersect at $\eta = 0$.

The influence of the thermophoresis parameter Nt and Brownian motion parameter Nb on the nanoparticles concentration φ are given in figures (9) and (10), respectively. Therefore, in these figures, equation (18) is evaluated versus the non-dimensional coordinate η . It is also noted from these figures that the nanoparticles concentration φ increases with the increase of Nt , whereas it decreases as Nb increases. It is also noted that the nanoparticles concentration for different values of Nt becomes larger with increasing η and reaches a maximum value (at a finite value of η : $\eta = \eta_0$) after which it decreases. The following clarifies the thermophoresis effect on nanoparticles concentration, namely, the result in Fig (9). Thermophoresis or thermo-migration is an original sin that occurs in a blend of transported particles, where the different particle types display different echoes to the nanoparticles concentration gradient force. So, the thermophoresis parameter causes an increment of the nanoparticles concentration. Fig. (11) shows the variation of the concentration of nanoparticles φ with η for several values of the parameter of thermal radiation R . It is seen from Fig.(11), that the nanoparticles concentration decreases with the increase of R , in the interval $\eta \in [0, 2.9]$; otherwise, it increases by increasing R . So, the manner of φ in the interval $\eta \in [2.9, \infty)$ is counteractive to its manner in the interval $\eta \in [0, 2.9]$. The other figures have the effect, so they are excluded here to save space. A dual role phenomena has obtained in figure (12) by increase the values of β_2 against the values of concentration distribution.

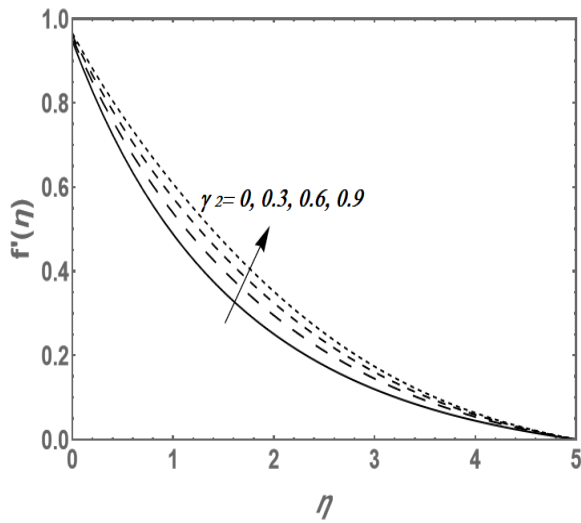


Fig. 1: Velocity ($f'(\eta)$) distribution versus γ_2 .

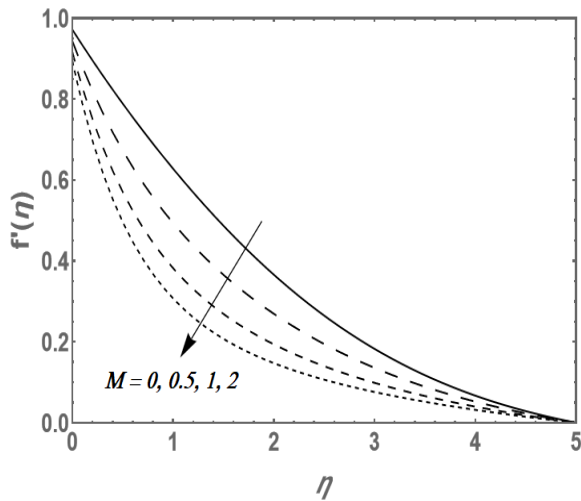


Fig. 2: Velocity ($f'(\eta)$) distribution versus M .

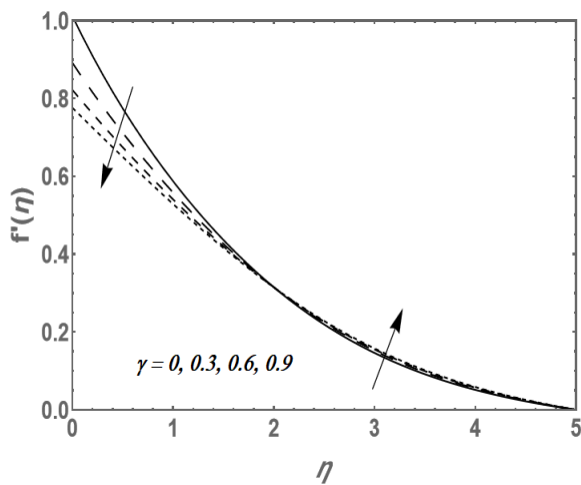


Fig. 3: Velocity ($f'(\eta)$) distribution versus γ .

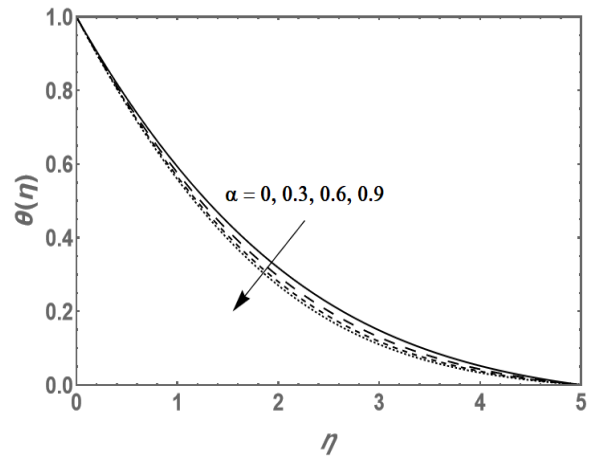


Fig. 4: Microrotation ($g(\eta)$) distribution versus α .

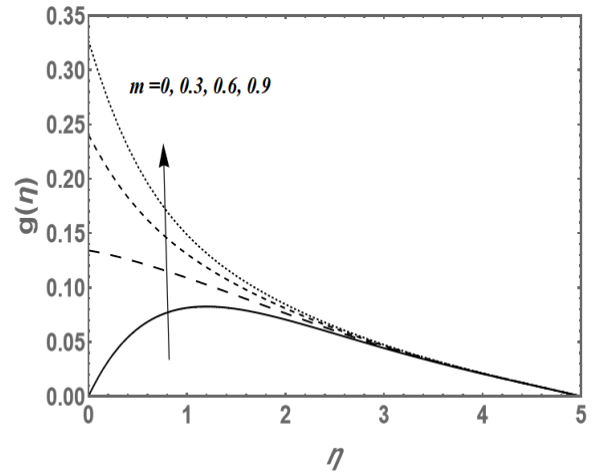


Fig. 5: Microrotation ($g(\eta)$) distribution versus m .

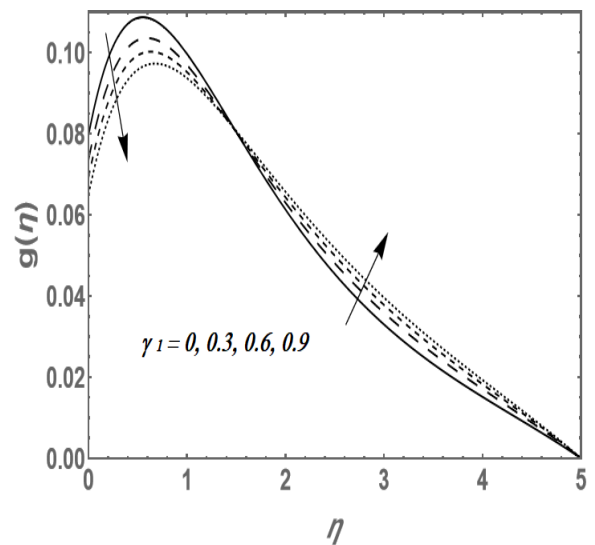


Fig. 6: Microrotation ($g(\eta)$) distribution versus γ_1 .

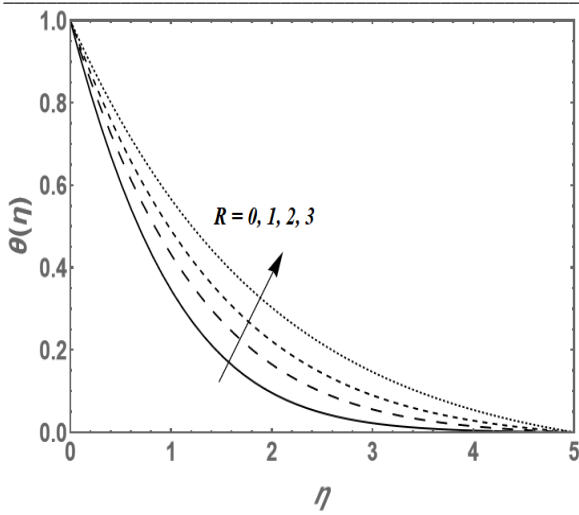


Fig. 7: Temperature ($\theta(\eta)$) distribution versus R .

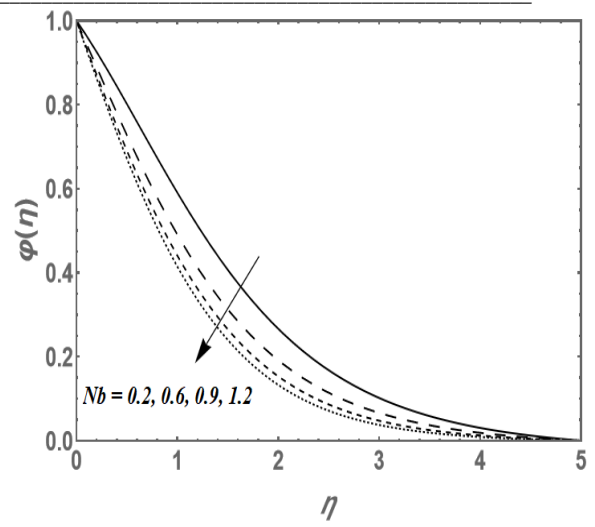


Fig. 10: Concentration ($\varphi(\eta)$) distribution versus N_b .

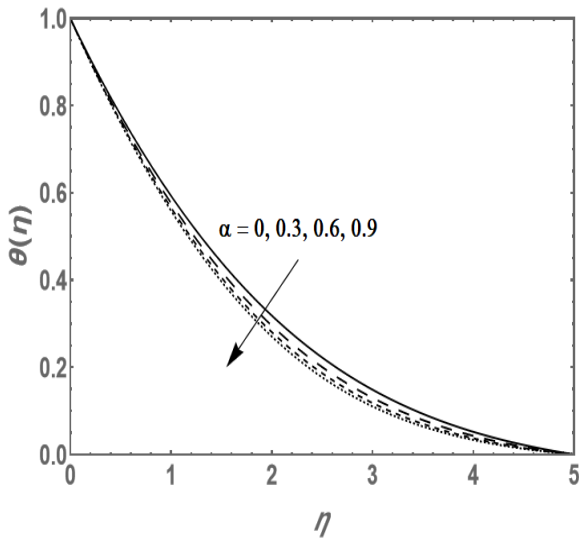


Fig. 8: Temperature ($\theta(\eta)$) distribution versus α .

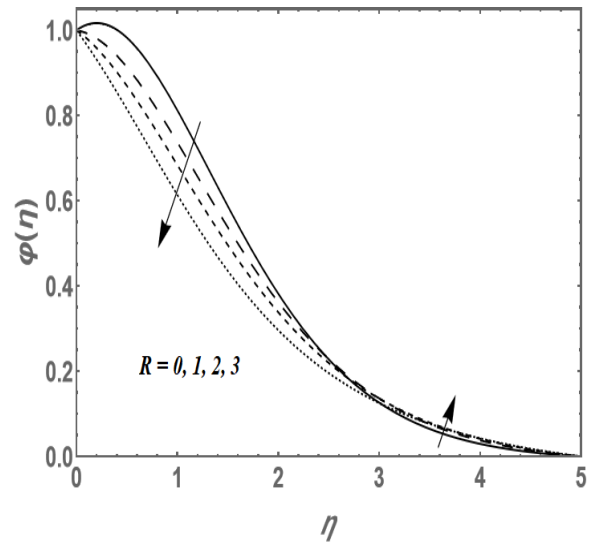


Fig. 11: Concentration $\varphi(\eta)$ distribution versus R .

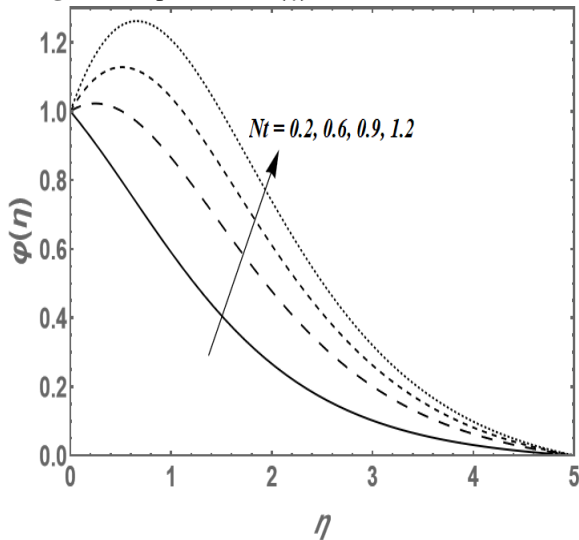


Fig. 9: Temperature ($\varphi(\eta)$) distribution versus M

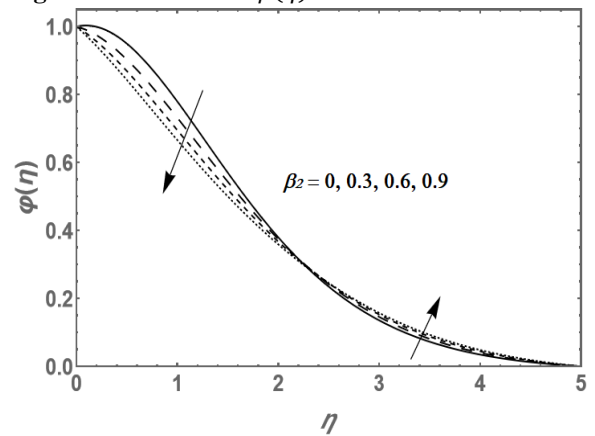


Fig. 12: Concentration $\varphi(\eta)$ distribution versus β_2 .

5 Conclusion

This analysis scrutinizes a new theoretical suggestion for a mechanism of electrical and thermal conductivity on the boundary layer flow of micro polar fluid with a stretching sheet. Multiple/microstructural slips, radiation, and nanoparticle are also studied. The fluid model is designated arithmetically by a system of ODEs. A computer program use the *Mathematica 13.1.1* is prepared with a shooting method to obtain the results/sketches. Investigations of these effects together are very useful due to their important vital applications in various scientific fields, especially in medicine and medical industries, such as endoscopes, respirators, and diverse medical implementations, as nanoparticles can be utilized in the remedy of cancer tumors. Moreover, this study investigates the influence of an endoscope on the unsteady incompressible flow which plays a very important role in medical diagnosis due to its wild clinical applications in determining the reasons behind many diseases in the human organs. For example, the motion of gastric juice when an endoscope is inserted through a small intestine. In addition, it is expected that the current effort will help a vehicle for considerate additional multifaceted difficulties in industry, engineering [41-74]. The main/essential notes are itemized as follows:

- The fluid temperature gets better at the higher values of thermophoresis (N_t) number in the case of ($\beta_1, \beta_2 = 0.3$).
- Mass Grashof (γ_2) number and Magnetic (M) parameter have a dual role phenomenon on the fluid-particle concentration.
- Thermophoresis (N_t) number has a clear/non-wobbling behavior on the fluid temperature and concentration.
- γ has two contradicting effects on the fluid velocity distribution.
- Micro rotation distribution has a stable behavior in absence of micro polar fluid.
- Results indicate that the concentration-dependent electrical and thermal conductivity gets the fluid particle distributions better.

Data availability statements

The authors states that all the files are provided in the paper no hidden file is required however if journal required any further data from us we will provide and the the corresponding author is responsible to provide to the journal.

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